



*From DC current heating to 3D
reconnection - the modes of magnetic
energy conversion in chromosphere and
corona - Numerical simulations based on
Hinode observations*

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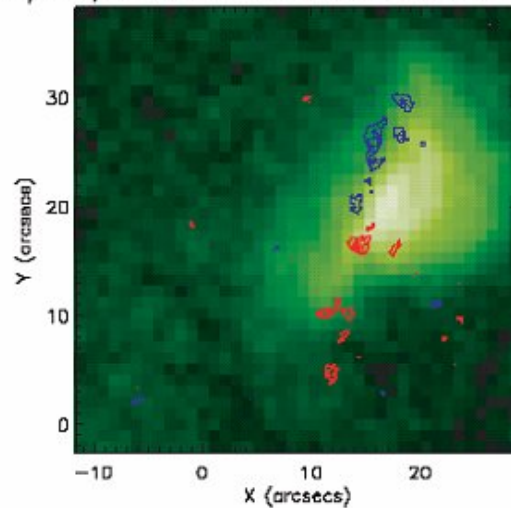
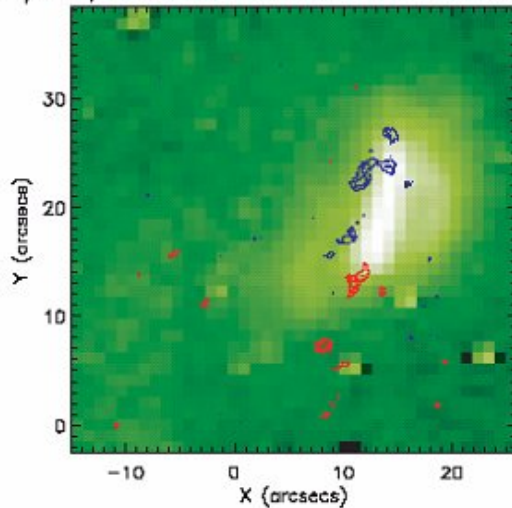
collaboration:

J. C. Santos, S. Jevadi (IMPRS Lindau)

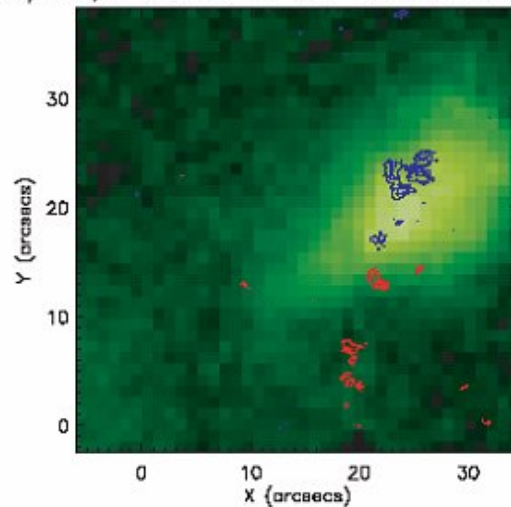
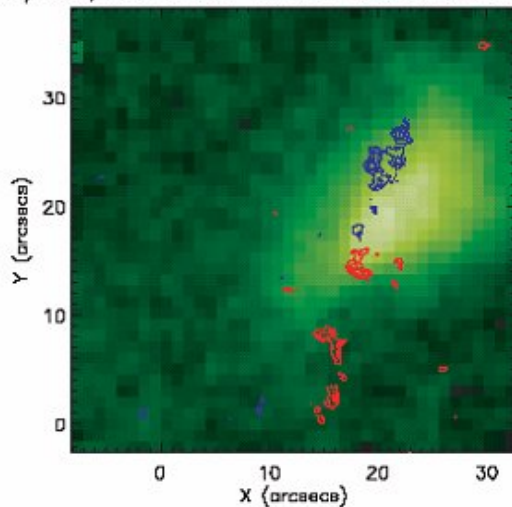
XRT and SOT observations



JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:00:44.397 UT JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:21:42.816 UT

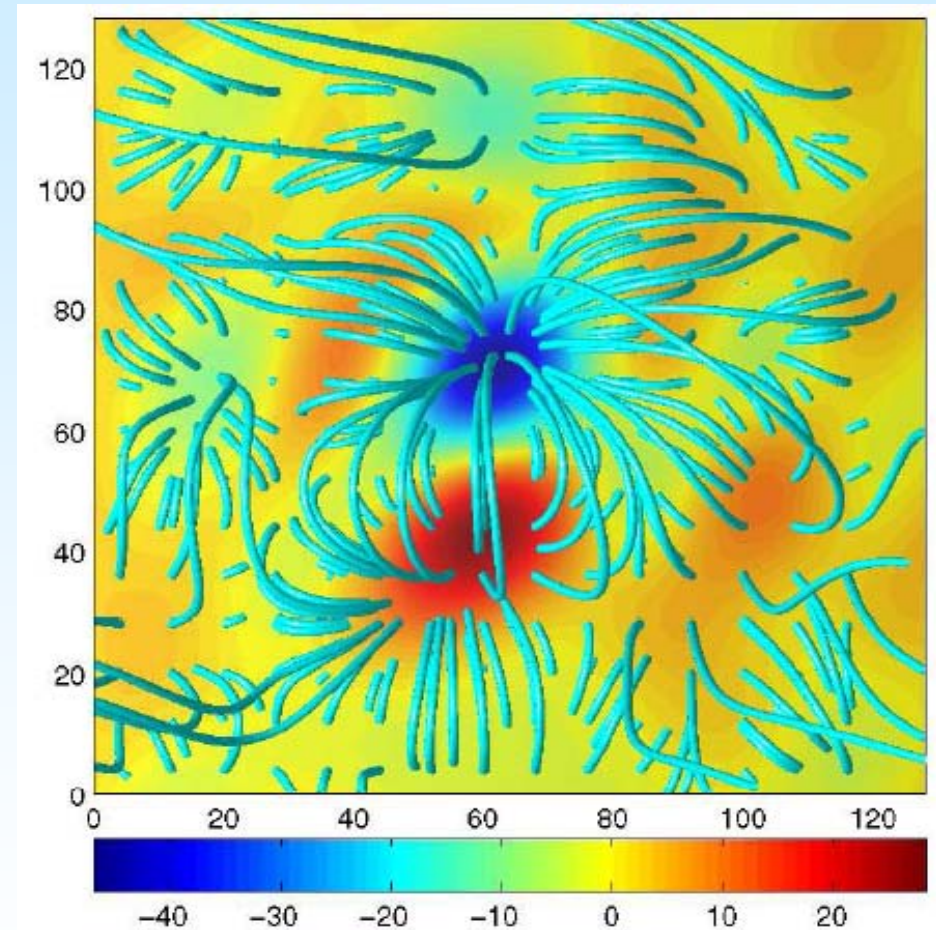
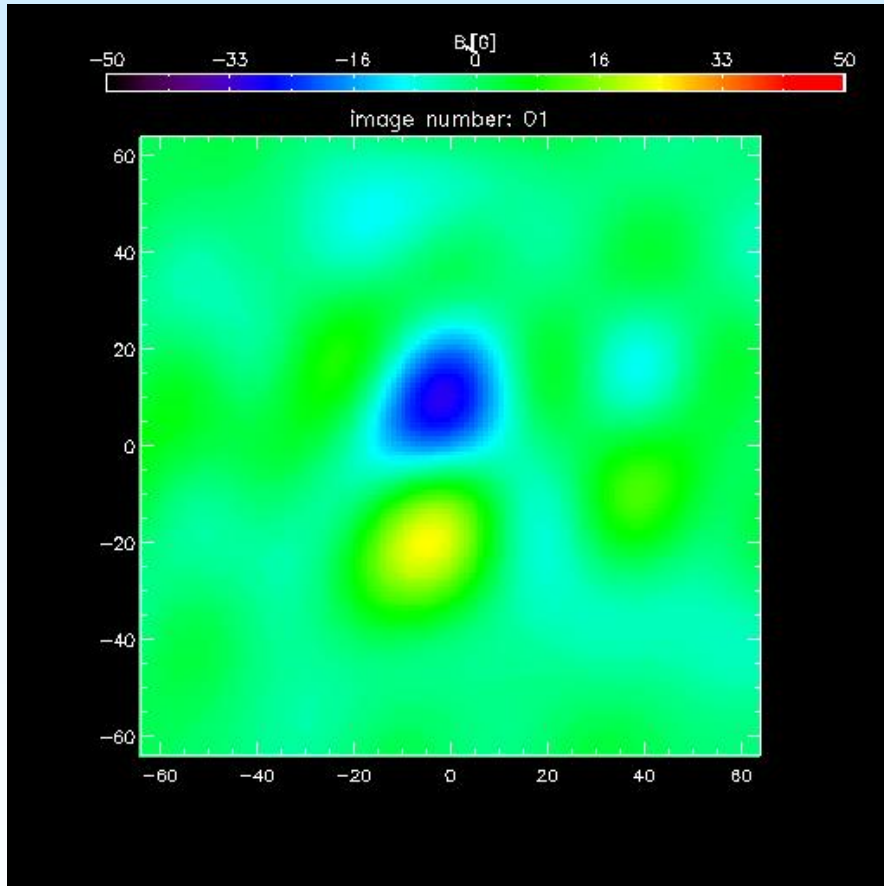


JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:43:03.241 UT JAXA/ISAS, SIRIUS XRT 19-Dec-2006 23:56:50.014 UT



XRT- X-ray
images taken
between
23:00 UT and
23:56 UT on
December
12, 2006
as well as SOT
photospheric
B-fields
from
[Kotoku et al.,
2007]
-> What causes
the brightening?

LOS photospheric B-field around the observed Xray-BP (23:00 UT)



Basic 8 spatial modes of the photospheric B-fields (LOS)

J. Büchner Modes of energy conversion

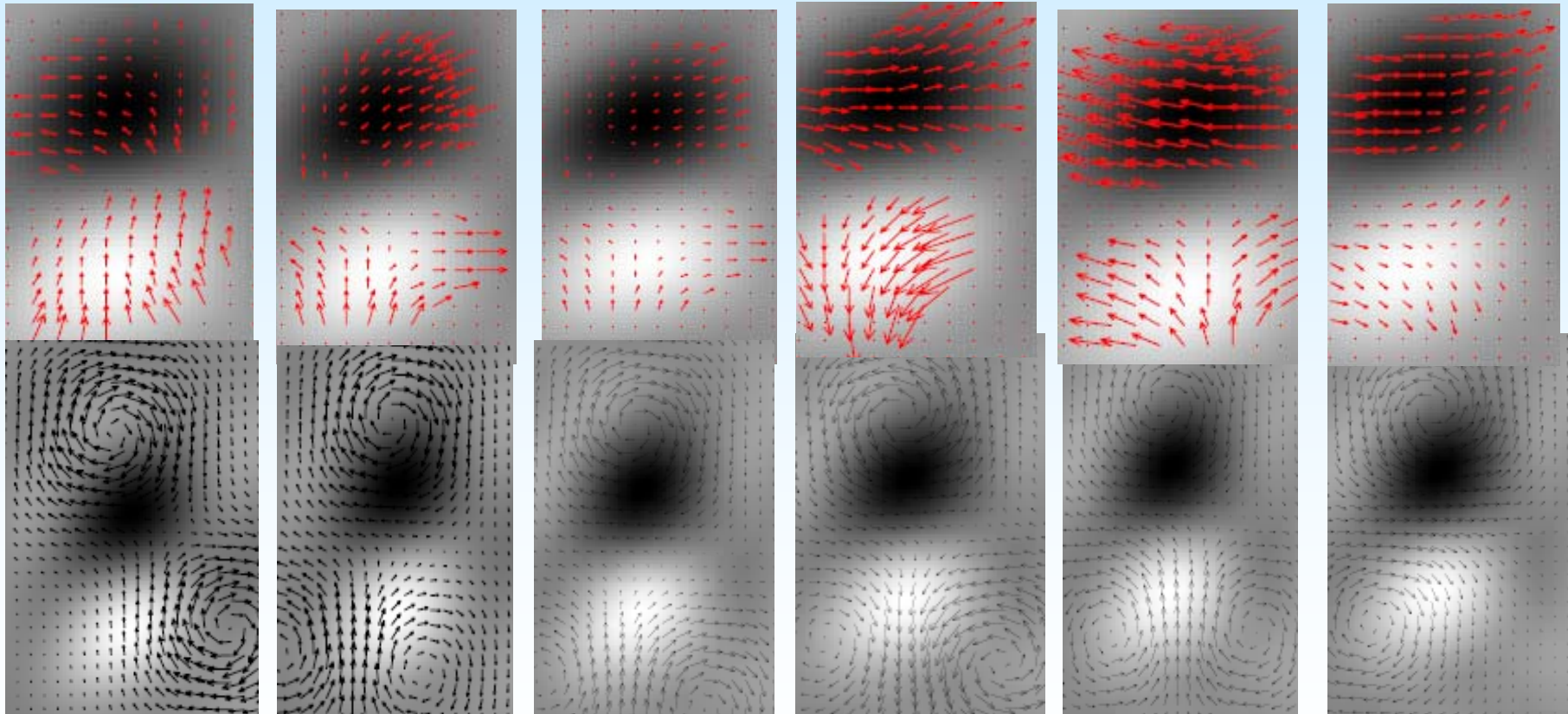
The potential B field shows Strong connection changes

2nd Hinode Science Meeting, Boulder, 30.9.08

LCT -> Patterns of photospheric plasma motion change with time

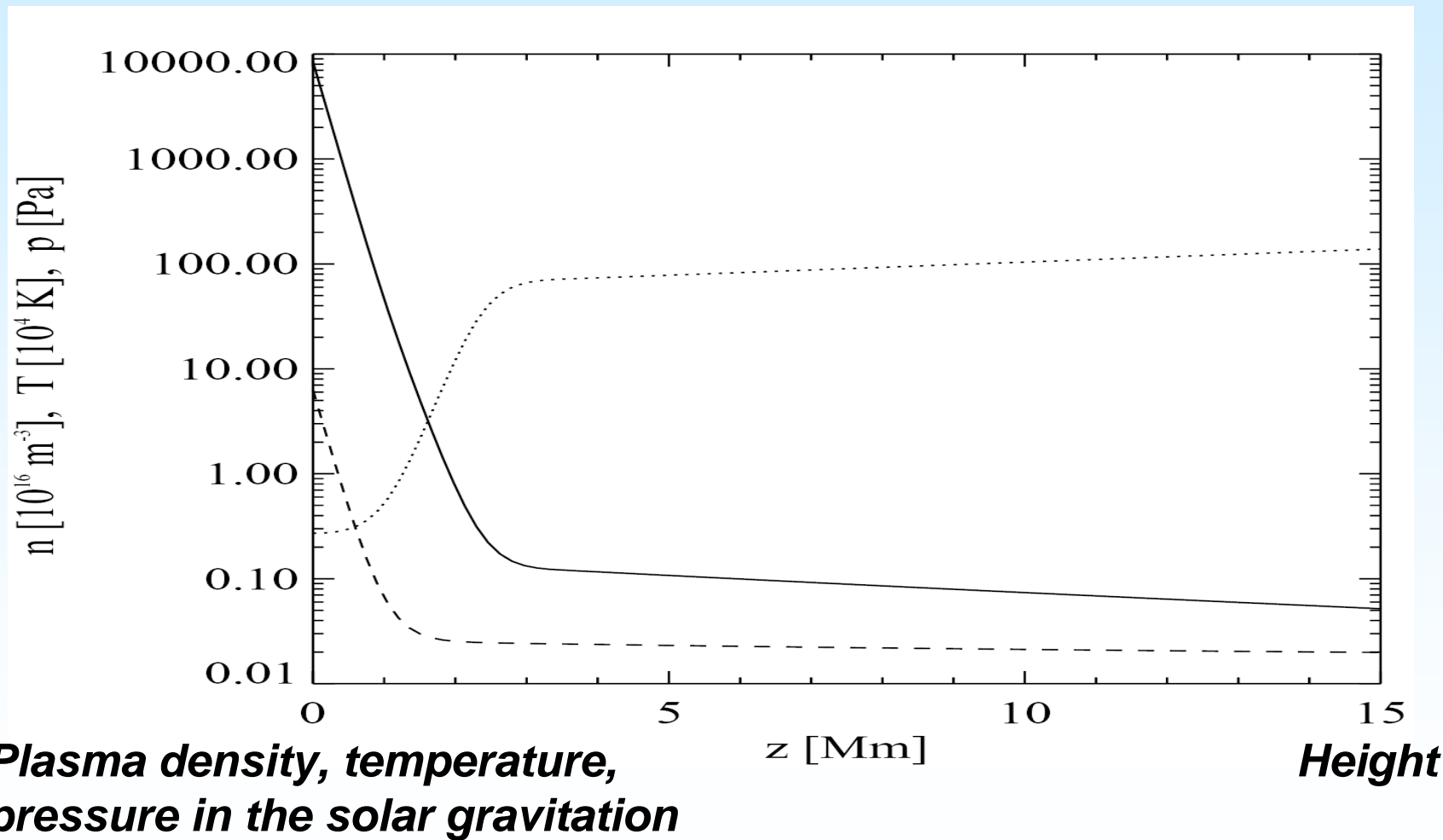


23:02 UT	23:07 UT	23:12 UT	23:17 UT	23:22 UT	23:27 UT
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Corresponding boundary conditions for simulation

Equilibrium plasma conditions: density and temperature stratification



Set of RMHD equations, solved



$$\begin{aligned}\frac{\partial \rho}{\partial t} &= -\vec{\nabla} \cdot \rho \vec{u} - \mu(\rho) && \text{(subscript 0 for neutrals)} \\ \frac{\partial \rho \vec{u}}{\partial t} &= -\vec{\nabla} \cdot \rho \vec{u} \vec{u} - \vec{\nabla} p + \vec{j} \times \vec{B} - \nu \rho (\vec{u} - \vec{u}_0) \\ &= -\vec{\nabla} \cdot \left[\rho \vec{u} \vec{u} + \left(p + \frac{B^2}{2\mu_0} \right) \mathbf{1} - \frac{\vec{B} \vec{B}}{\mu_0} \right] - \nu \rho (\vec{u} - \vec{u}_0) \\ \frac{\partial \vec{B}}{\partial t} &= \vec{\nabla} \times (\vec{u} \times \vec{B} - \eta \vec{j}) \\ \frac{\partial p}{\partial t} &= -\vec{\nabla} \cdot p \vec{u} - (\gamma - 1) p \vec{\nabla} \cdot \vec{u} + (\gamma - 1) \eta j^2 - \kappa n k_B (T - T_0)\end{aligned}$$

$$\begin{aligned}\vec{E} &= -\vec{u} \times \vec{B} + \eta \vec{j} \\ \vec{\nabla} \times \vec{B} &= \mu_0 \vec{j}\end{aligned}$$

$$p = 2n\kappa_B T$$

$$\hat{\eta}^* = \begin{cases} \eta_a^* \left(\frac{|v_{\text{dr}}|}{v_{\text{thr}}} - 1 \right); & |v_{\text{dr}}| \geq v_{\text{thr}}, \\ 0; & |v_{\text{dr}}| < v_{\text{thr}}. \end{cases}$$

RMHD simulation results

- Simulated heating is maximum at BPs position
 - The heating is robust, **does not depend** on the details of the photospheric plasma motion.
 - **Most important for heating are geometrical features of the magnetic field, in particular strong changes in the B-field connectivity**
 - **A consequence of connectivity changes if footpoint motion applies: current generation**
- To understand current dissipation:**
-> kinetic theories have to be applied

Current dissipation

If the resistivity is parametrized by an „effective collision frequency“ as

$$\eta = \frac{\nu}{\epsilon_0 \omega_{pe}^2}$$

then in the **chromosphere binary particle collisions** dissipate currents [Spitzer–Härm]:

$$\nu_{coll} \cong \frac{\omega_{pe}}{n \lambda_D^3}$$

However, in the **transition region and corona -> binary collisions become inefficient->micro-turbulent dissipation:**

Threshold: current carrier velocity or gradients > critical Quasi-linear estimate [Sagdeev et al.]

$$\nu_S \cong 0.01 \omega_{pe} \frac{T_e}{T_i} \frac{u_D}{v_{te}}$$

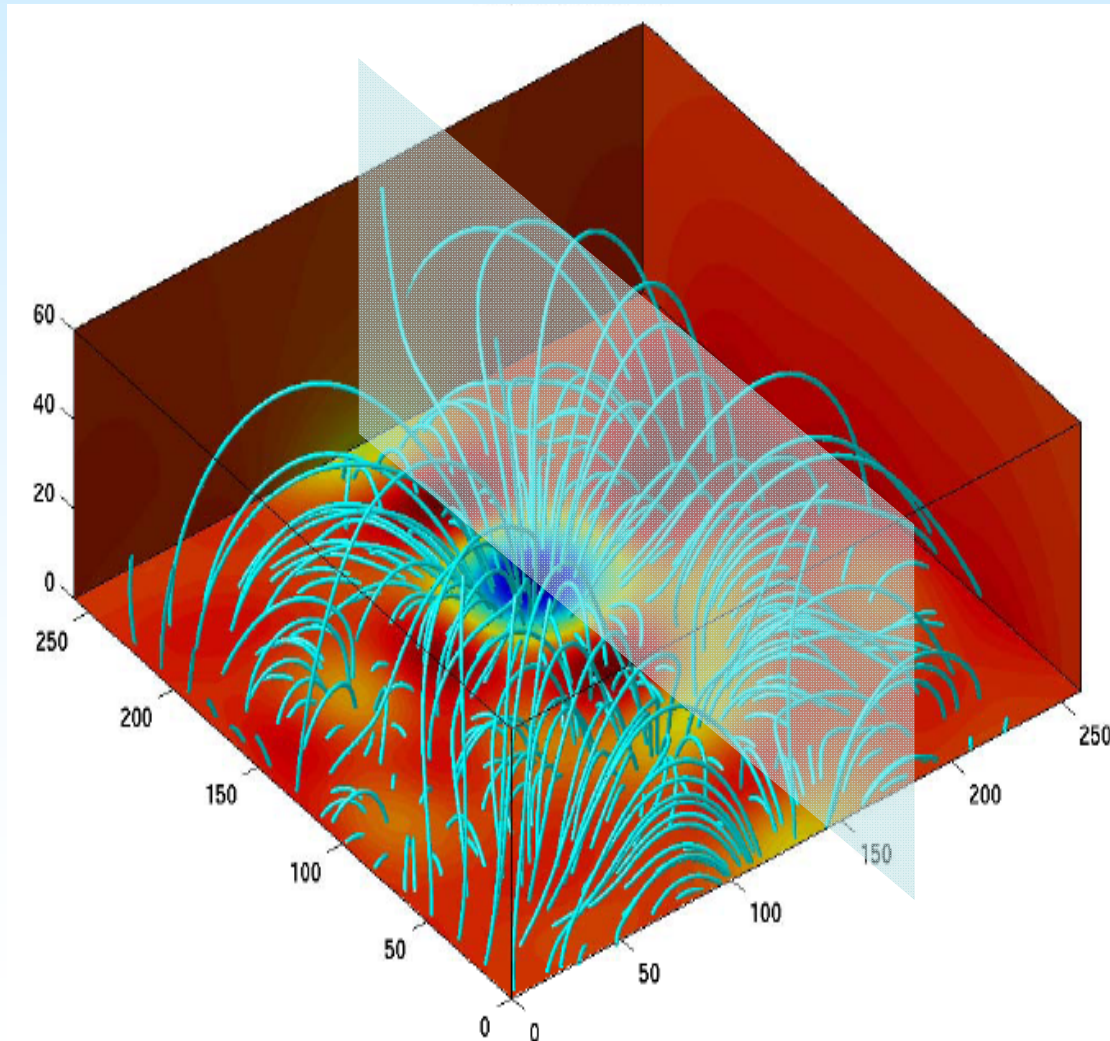
Nonlinear->kinetic simulations needed for solar conditions [Büchner & Elkina]

- higher beta -> IA / double layers
- lower beta -> LH turbulence

$$\nu_c \approx \omega_{pi} / 2\pi$$

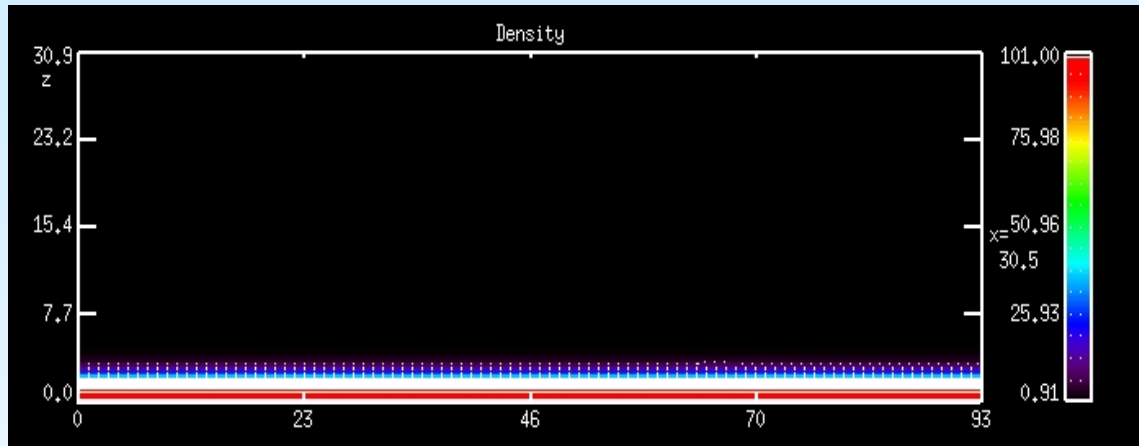
$$\nu_c \approx \Omega_{LH}$$

Results of simulations diagnosed in a vertical plane through the B_{max}

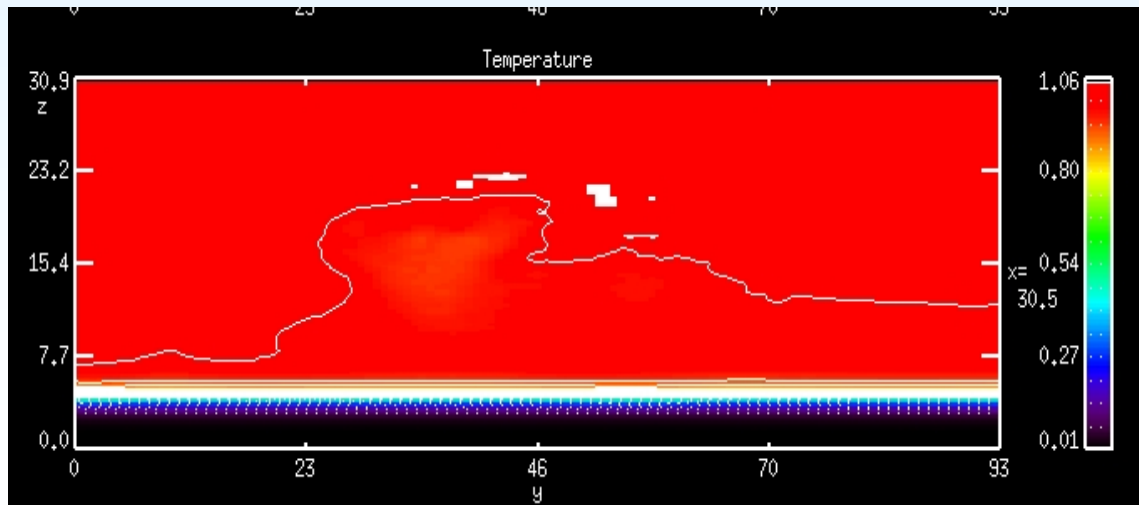


Initial situation

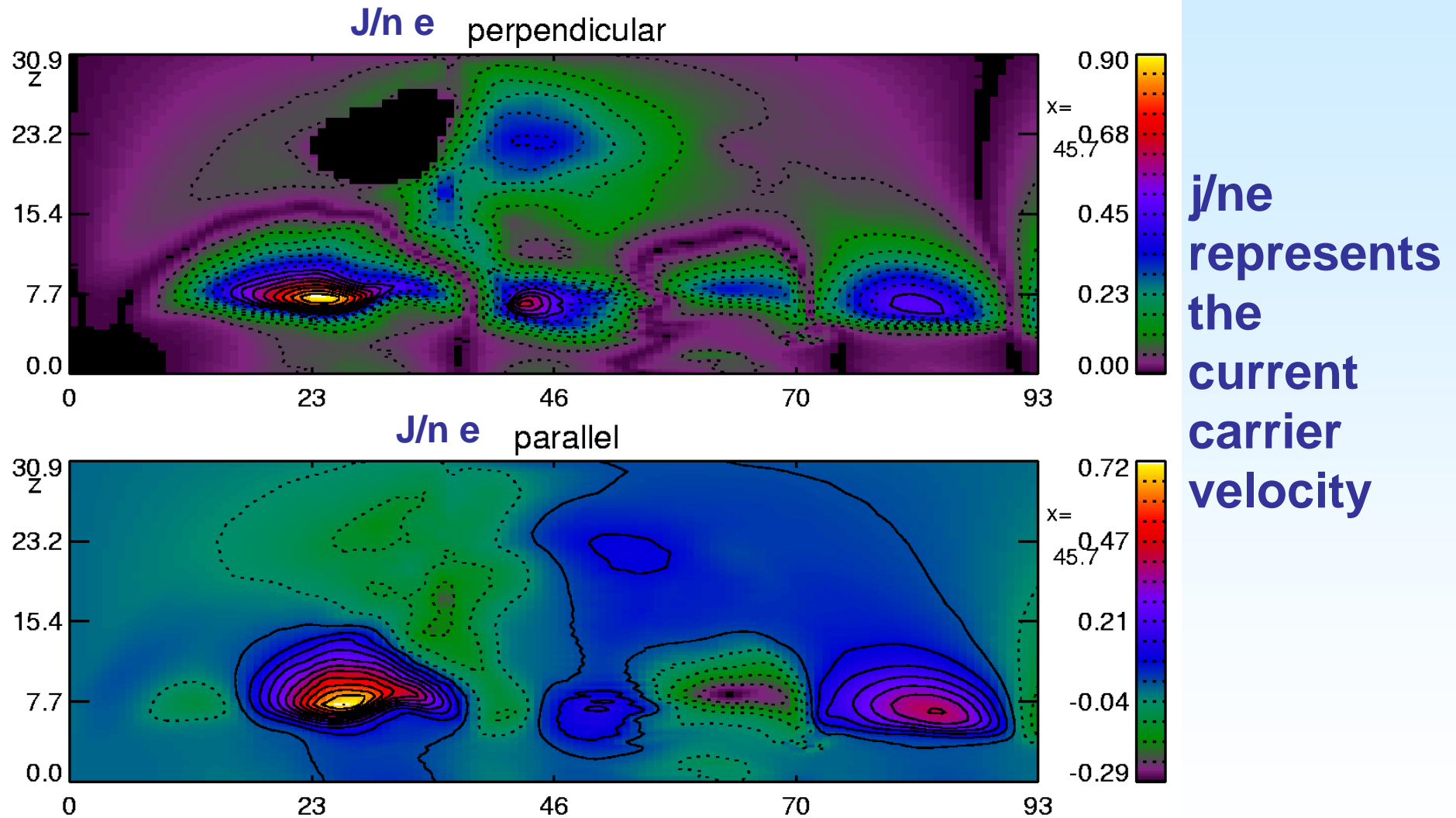
Density (n)



**Temperature
(T)
(Red = 10^6 K)**

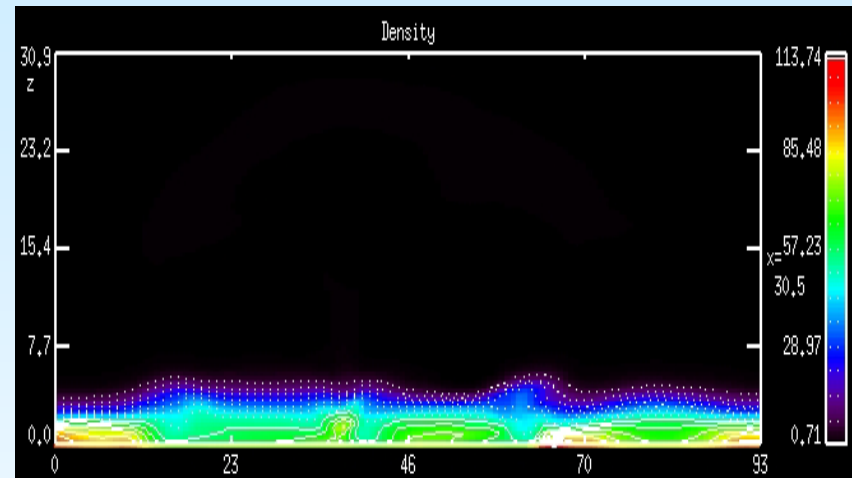
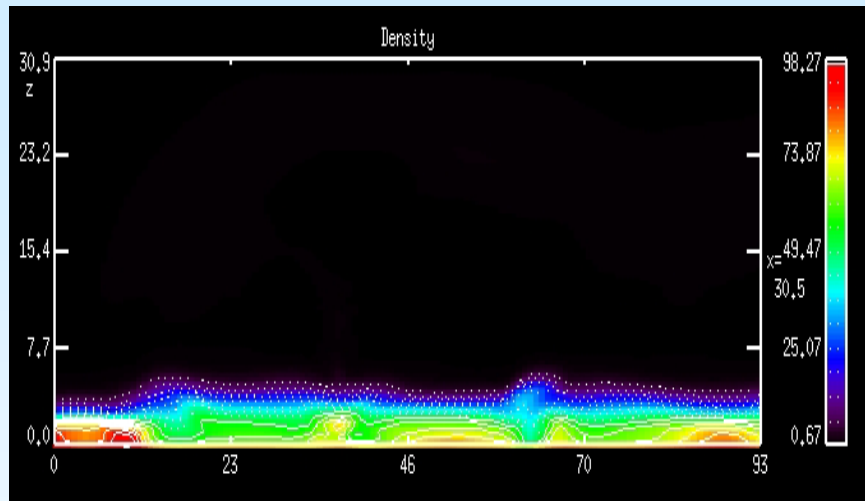


Turbulent current dissipation

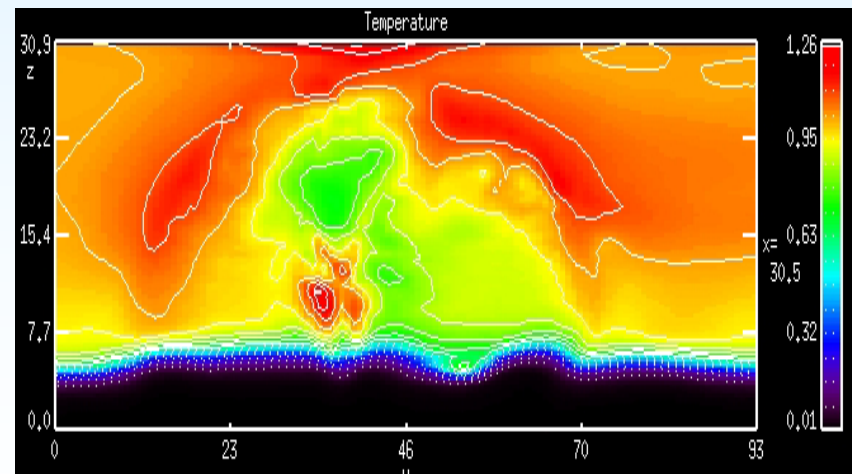
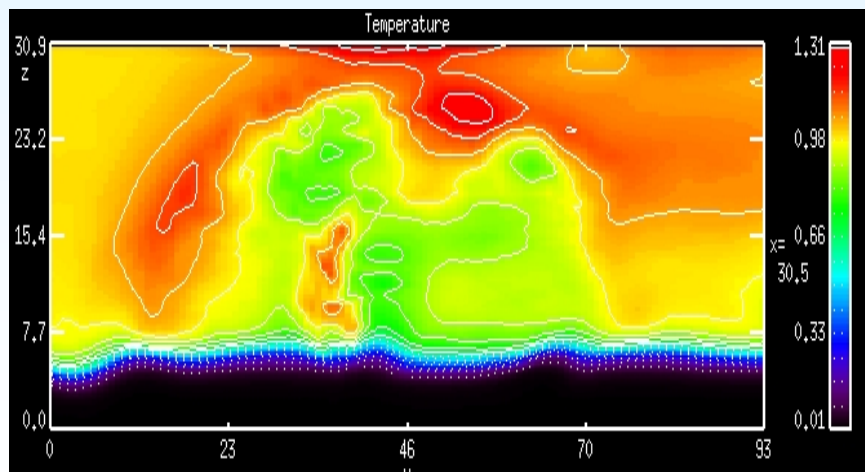


Results of 5-minute runs

23:02 – 23:07 & 23:07 – 23:12



n

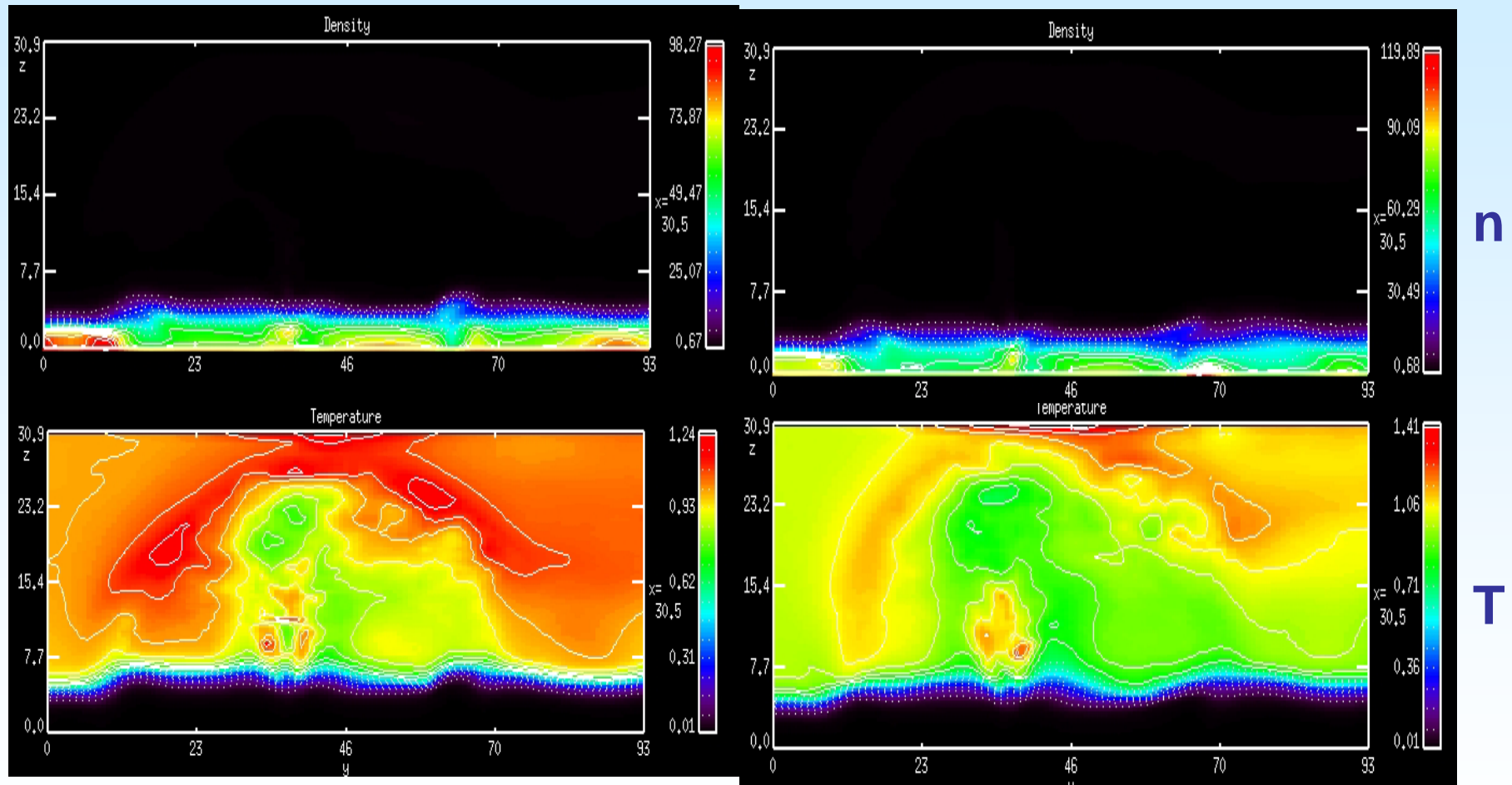


T

(Red = $1.3 \cdot 10^6$ K)

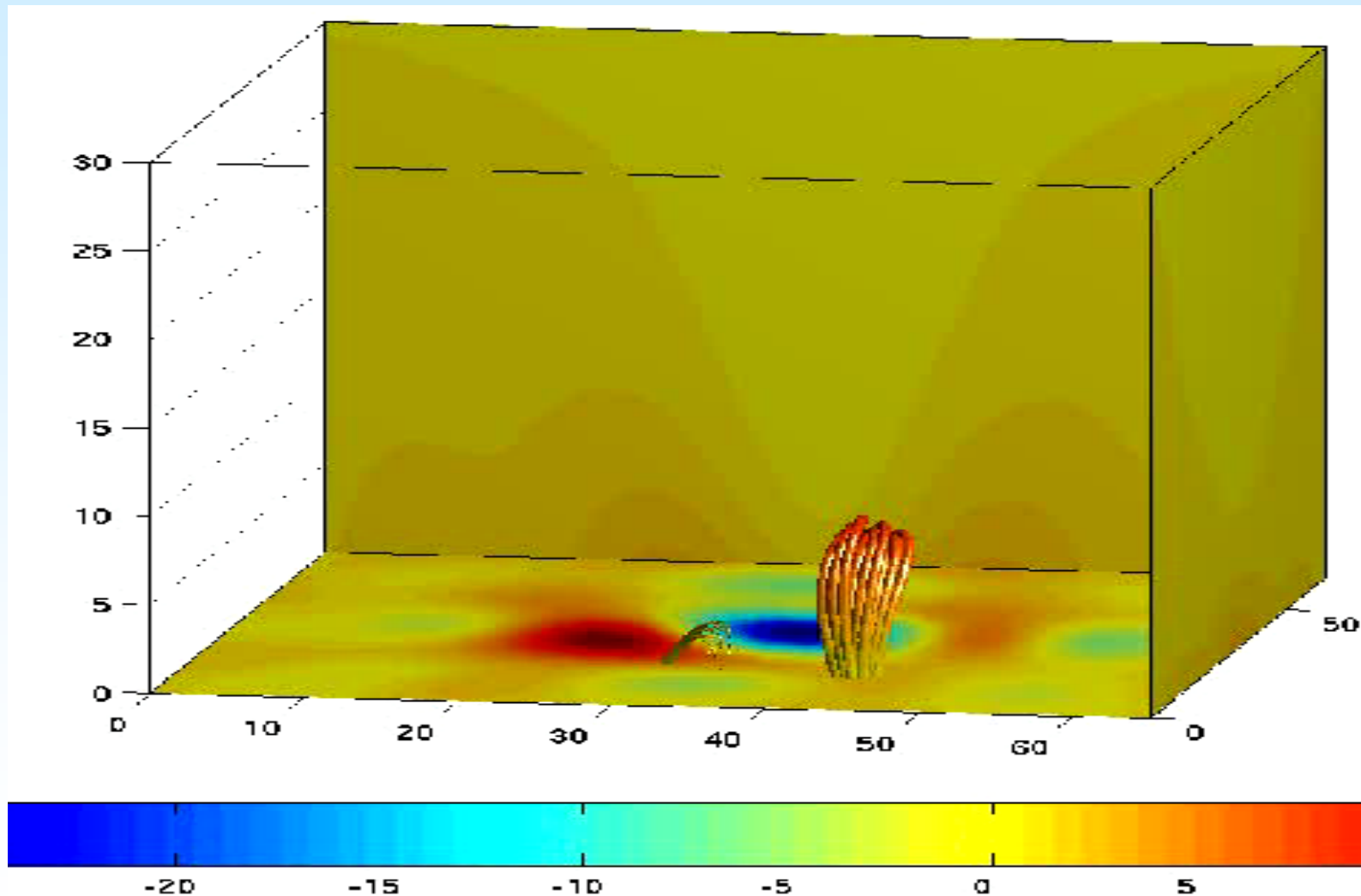
Results of 5-minute runs

23:12 – 23:17 & 23:17 – 23:22



(Red = $1.4 \cdot 10^6$ K)

Above the chromosphere: 3D reconnection („slipping“)



Conclusions



- An observed Hinode X-ray Bright Point location
 - coincides well with the sites of strong B-field **connectivity changes**, and
 - **do not depend** much on the exact form of the **plasma motion** in the photosphere below
- Strong B-field connectivity changes + some plasma motion across the footpoints
 - > **currents, both parallel and perpendicular to B**
- While in the chromosphere binary collisions dominate the current dissipation, causing Ohmic heating,
- In the transition region / lower corona main dissipation microturbulence-induced 3D „slipping“ reconnection

Outline



- Unprecedented high resolution Hinode observations
- But: Still open, what causes the formation of X-ray and EUV Bright Points
 - coincides well with the sites of strong B-field **connectivity changes**, and
 - **do not depend** much on the exact form of the **plasma motion** in the photosphere below
- Strong B-field connectivity changes + some plasma motion across the footpoints
 - > **currents, both parallel and perpendicular to B**
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- In the transition region / lower corona main dissipation microturbulence-induced 3D „slipping“ reconnection